



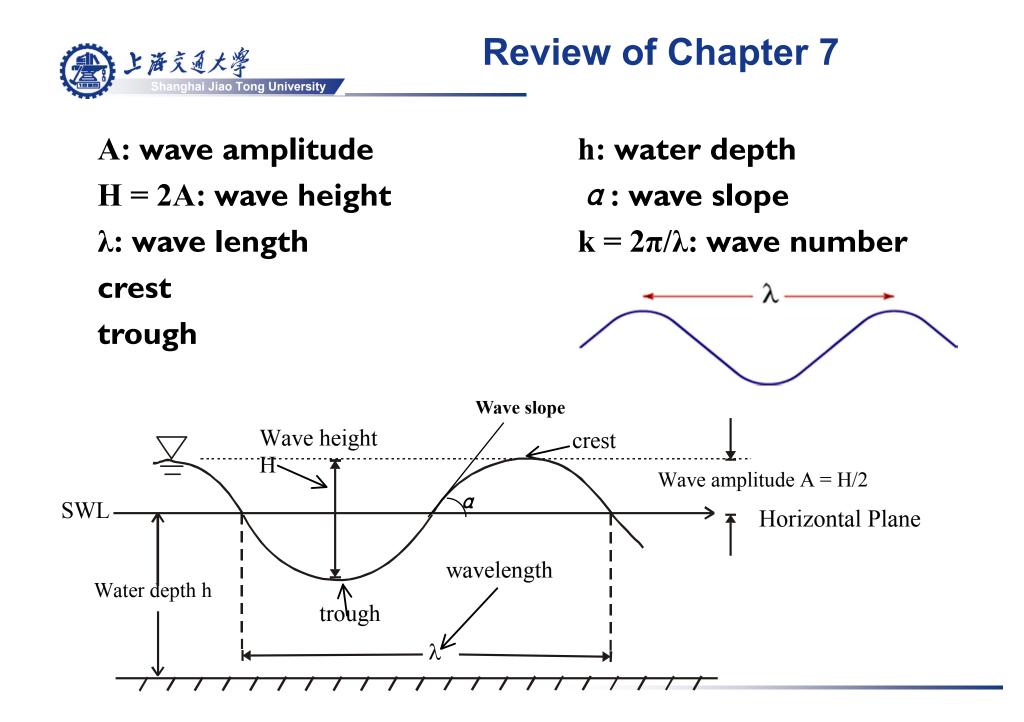
Introduction to Marine Hydrodynamics (NA235)

Department of Naval Architecture and Ocean Engineering School of Naval Architecture, Ocean & Civil Engineering

Shanghai Jiao Tong University

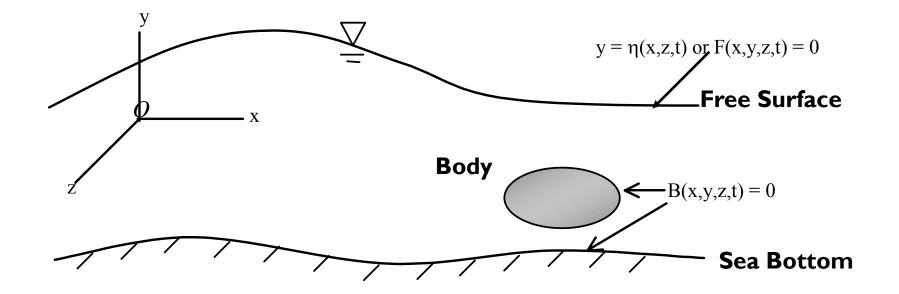


Chapter 7 Water Waves





Governing equation of water waves



A summary of governing equations of water waves.

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Field equations: in $y \leq \eta(x, z, t)$ **Laplace's Eq.** $\nabla^2 \phi = 0$ Dynamic cond. $\frac{\partial \phi}{\partial t} + \frac{\left|\nabla \phi\right|^2}{2} + \frac{p - p_a}{\rho} + gy = 0$ $\frac{\partial \phi}{\partial t} \to 0, \quad \mathbf{V} = \nabla \phi \to 0, \quad p = p_a - \rho g y$ Far field cond. $\begin{cases} \phi(x,\eta,z,0) = f(x,z) & \text{on } y = \eta \\ \eta(x,z,0) = g(x,z) \end{cases}$ Initial cond.

Boundary conditions:

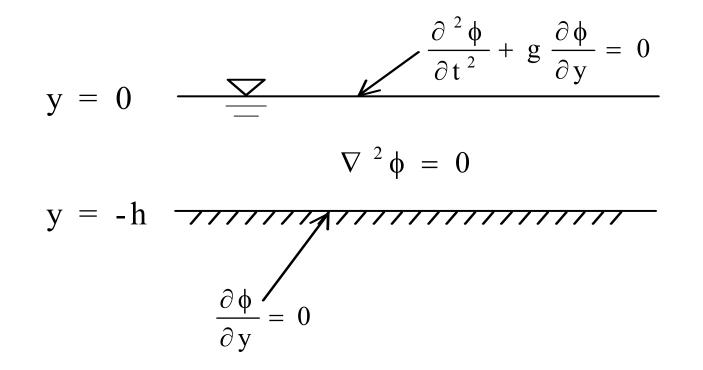
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On body surf., B	$\frac{\partial \phi}{\partial \mathbf{n}} = \mathbf{U}_n \text{or} \frac{\partial B}{\partial t} + (\nabla \phi \cdot \nabla) B = 0$
On the bottom, $y = -h$	$\frac{\partial \phi}{\partial y} = 0$
On free surface, $y = \eta$ (kinematic cond.)	$\frac{\partial \phi}{\partial y} = \frac{\partial \eta}{\partial t} + \frac{\partial \phi}{\partial x} \frac{\partial \eta}{\partial x} + \frac{\partial \phi}{\partial z} \frac{\partial \eta}{\partial z}$
On free surface, $y = \eta$ (dynamic cond.)	$\frac{\partial \phi}{\partial t} + \frac{\left \nabla \phi\right ^2}{2} + g\eta = 0$



Governing equations of Airy wave



Linearized governing equations of Airy waves are as follows.

<u>Field equations</u> in $y \le \eta(x, z, t)$

I. Laplace's Eq. $\nabla^2 \phi = 0$

2. Dynamic cond.
$$p - p_a = -\rho \frac{\partial \phi}{\partial t} - \rho g y$$

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4. Initial cond.

$$\frac{\partial \phi}{\partial t} \to 0, \quad \mathbf{V} = \nabla \phi \to 0, \quad p = p_a - \rho g y$$
$$\begin{cases} \phi(x, 0, z, 0) = f(x, z) \\ \eta(x, z, 0) = g(x, z) \end{cases}$$

Boundary conditions:

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5. On body surface, B

$$\frac{\partial \phi}{\partial \mathbf{n}} = \mathbf{U}_n \text{ or } \frac{\partial B}{\partial t} + (\nabla \phi \cdot \nabla) B = 0$$

6. On the bottom, y = -h

$$\frac{\partial \phi}{\partial y} = 0$$

7. On the mean free surface, y = 0

$$g \frac{\partial \phi}{\partial y} + \frac{\partial^2 \phi}{\partial t^2} = 0$$
$$\eta = -\frac{1}{g} \frac{\partial \phi}{\partial t}$$



Velocity potential function of Airy waves can be derived.

$$\phi = \frac{gA}{\omega} \sin\left(kx - \omega t\right) \frac{\cosh k\left(y + h\right)}{\cosh kh}$$

As the depth of water tends to infinity, $h \rightarrow \infty$, it becomes

$$\phi = \frac{gA}{\omega} \sin\left(kx - \omega t\right) e^{ky}$$



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$$\eta = -\frac{1}{g} \frac{\partial \phi}{\partial t} \bigg|_{y=0} = A \cos(kx - \omega t)$$

From the dynamic condition, pressure distribution of Airy wave is

$$p - p_a = -\rho \frac{\partial \phi}{\partial t} - \rho gy$$
$$= -\rho gy + \rho Ag \cos(kx - \omega t) \frac{\cosh k(y+h)}{\cosh kh}$$



Velocity field of Airy wave

$$u = \frac{\partial \phi}{\partial x} = A\omega \frac{\cosh k (y+h)}{\sinh kh} \cos (kx - \omega t)$$
$$v = \frac{\partial \phi}{\partial y} = A\omega \frac{\sinh k (y+h)}{\sinh kh} \sin (kx - \omega t)$$



Dispersion relation of Airy wave

$$\omega^{2} = gk \frac{\sinh kh}{\cosh kh} = gk \tanh kh$$
$$c^{2} = \frac{\omega^{2}}{k^{2}} = \frac{g\lambda}{2\pi} \tanh\left(\frac{2\pi h}{\lambda}\right)$$



Water particle orbit of Airy wave

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$$\frac{\left(x_{p}-\overline{x}\right)^{2}}{a^{2}} + \frac{\left(y_{p}-\overline{y}\right)^{2}}{b^{2}} = 1$$
$$a = A \frac{\cosh k \left(\overline{y}+h\right)}{\sinh kh},$$
$$b = A \frac{\sinh k \left(\overline{y}+h\right)}{\sinh kh}$$



Characteristics of Airy waves

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	deep water	shallow water
	<i>kh</i> > 3	$kh \ll 1$
dispersion	$\omega^2 = gk$	$\omega^2 = ghk^2$
relation	$c_d^2 = \frac{g}{k} = \frac{g\lambda}{2\pi}$	$c_s^2 = gh$
velocity	$\frac{u}{U_0} \approx e^{ky}$	$\frac{u}{U_0} \approx 1$
	$\frac{v}{V_0} \approx e^{ky}$	$\frac{v}{V_0} \approx 1 + \frac{y}{h}$

Review of Chapter 7



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	deep water $kh > 3$	shallow water $kh << 1$
Water particle orbit	$a = b = A e^{k \overline{y}}$	$a = \frac{A}{kh}$ $b = A\left(1 + \frac{\overline{y}}{h}\right)$
Pressure field	$p - p_a = \rho g \left(\eta e^{k y} - y \right)$	$p - p_a = \rho g (\eta - y)$



Characteristics of Airy waves

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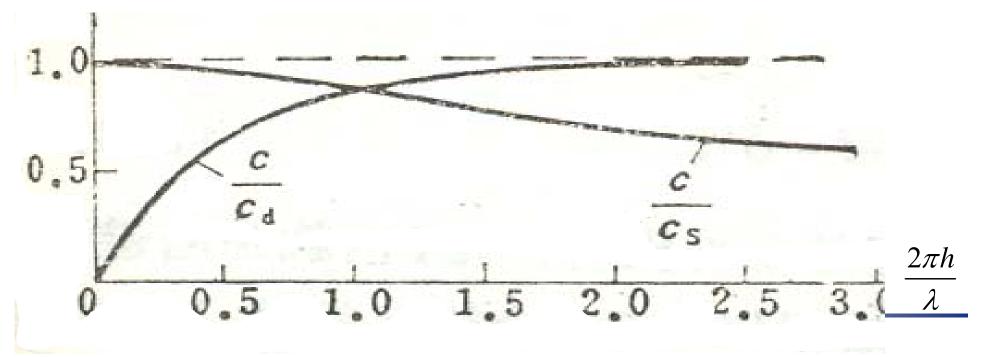
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	deep water	shallow water
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Water particle orbit	$a = b = A e^{k \overline{y}}$	$a = \frac{A}{kh}$ $b = A\left(1 + \frac{\overline{y}}{h}\right)$
Pressure field	$p - p_a = \rho g \left(\eta e^{k y} - y \right)$	$p - p_a = \rho g (\eta - y)$



For a fixed water depth, since $c_s = \sqrt{gh}$ is a constant, phase velocity c will decrease with wave length λ . That is, longer waves travel faster, and shorter waves travel slower.

For waves with fixed wave length, since $c_d = \sqrt{\frac{g\lambda}{2\pi}}$ is a constant, phase velocity c will increase with the water depth h. That is, deep water waves travel faster, and shallow water waves travel slower.





Linear plane progressive waves

$$\phi = \frac{gA}{\omega} \frac{\cosh k \left(y + h \right)}{\cosh kh} \sin \left(kx \cos \theta + kz \sin \theta - \omega t \right)$$

$$\eta = A\cos\left(kx\cos\theta + kz\sin\theta - \omega t\right)$$

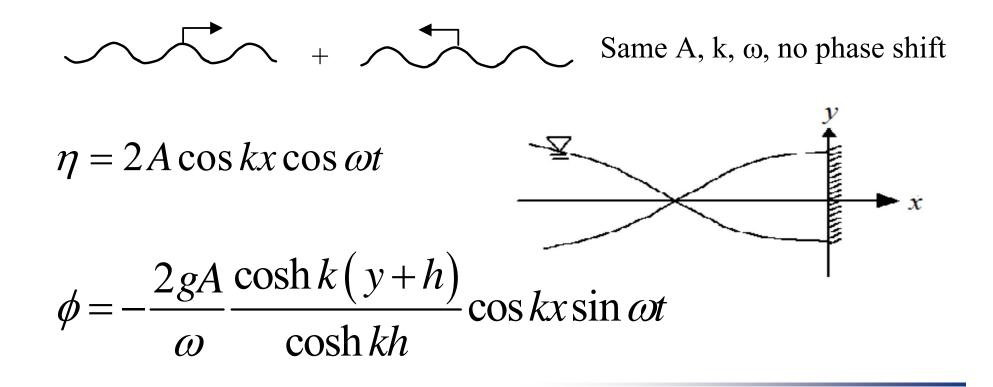
$$\omega^2 = gk \tanh kh$$

$$k_x = k \cos \theta, \quad k_z = k \sin \theta, \quad k = \sqrt{k_x^2 + k_z^2}$$



Linear superposition law

Standing Waves ---- a superposition of two Airy waves with same parameters, but propagating at opposite directions.



Oblique standing waves -- superposition of two plane progressive waves propagating not collinear.

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$$\eta_{T} = 2A\cos\left(kx\cos\theta\right)\cos\left(kz\sin\theta - \omega t\right)$$

$$\phi = -\frac{2gA}{\omega}\frac{\cosh k\left(y+h\right)}{\cosh kh}\cos\left(kx\cos\theta\right)\sin\left(kz\sin\theta - \omega t\right)$$

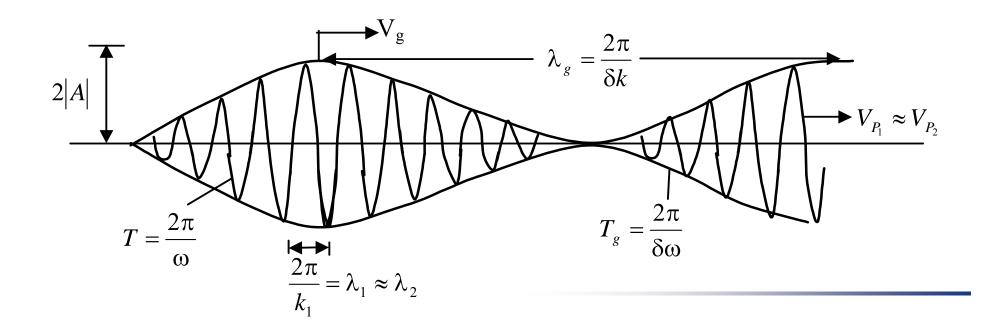
$$\eta_{T}$$

Wave Group: a superposition of different progressive waves with similar frequencies propagating along the same direction.

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$$\eta_T = \eta_1 + \eta_2 = \operatorname{Re}\left\{Ae^{i(k_1x - \omega_1t)} \left[1 + e^{i(\delta k x - \delta \omega t)}\right]\right\}$$





$$V_g = \frac{\delta\omega}{\delta k} = \frac{d\omega}{dk}$$

Group velocity of Airy waves (1/2 < C < 1)

$$V_{g} = \frac{d\omega}{dk} = \frac{\omega}{\frac{k}{V_{p}}} \underbrace{\frac{1}{2} \left(1 + \frac{2kh}{\sinh 2kh}\right)}_{C} = CV_{p}$$

For deep water, $C = \frac{V_{g}}{V_{p}} = \frac{1}{2}$; For shallow water, $C = \frac{V_{g}}{V_{p}} = 1$

Wave energy of Airy wave

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$$\overline{E} = \overline{PE}_{\text{wave}} + \overline{KE}_{\text{wave}} = \frac{\rho g A^2}{4} + \frac{\rho g A^2}{4} = \frac{\rho g A^2}{2}$$

Waver energy propagation equation

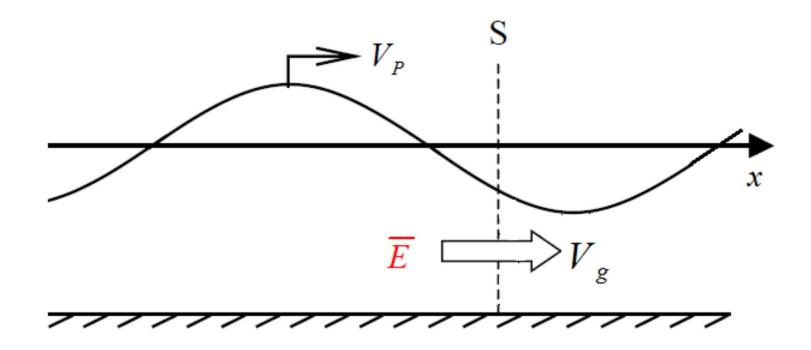
$$\frac{\partial \overline{E}}{\partial t} + \frac{\partial \left(V_g \overline{E} \right)}{\partial x} = 0$$



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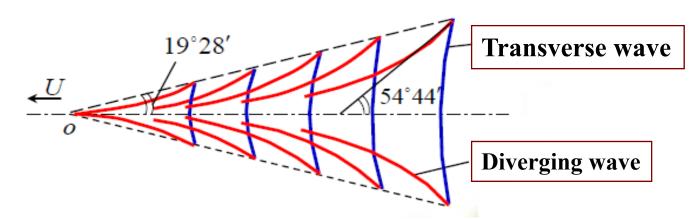
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wave energy propagates through vertical plane, perpendicular to the wave propagating direction, at group velocity.





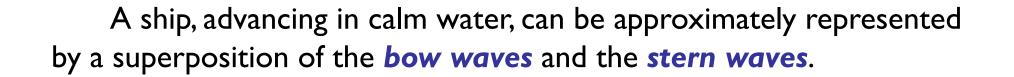
Kelvin waves

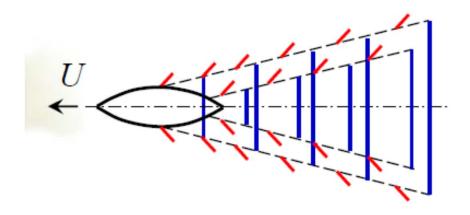


Two kind of waves in a moving point pressure generated wave systems 1. *Transverse waves*, nearly perpendicular to forward speed, of wave length $\lambda = 2\pi U^2/g$ 2. *Diverging waves*

Waves are confined to a sector of semi-angle 19° 28', known as *Kelvin angle*.







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Wave making resistance of a ship has relationship with the interference of bow wave and stern wave. Lower ship speed corresponds remarkable transverse wave, and higher speed corresponds remarkable diverging wave. Diverging waves disappear more rapidly than transverse waves.



Wave making resistance of a moving ship is equal to half of the average wave energy.

$$R_W = \frac{1}{2}\overline{E} = \frac{1}{4}\rho g A^2$$

Wave making resistance is proportional to the square of the wave amplitude, which is determined by the ship speed, hull form and the interference between bow waves and stern waves.



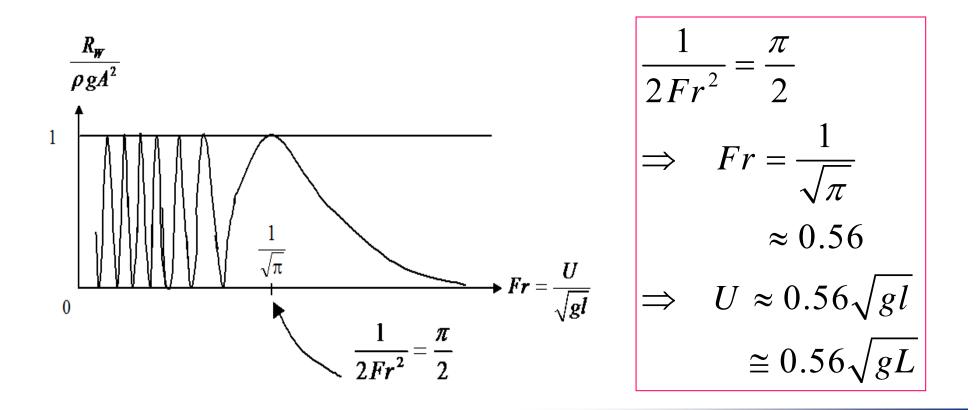
Waves behind a ship can be approximately looked as a superposition of a *bow wave* and a *stern wave*. Then, the wave resistance is expressed as

$$R_W = \rho g A^2 \sin^2 \left(\frac{1}{2Fr^2}\right)$$

For relatively lower speed, if the ship is well designed, the bow wave and the stern wave may cancel out with each other and wave resistance is small. But for higher speed ship, wave resistance is unavoidable, would not vanish.



If Froude number $F_r > 0.56$, wave resistance is unavoidable.





Chapter 8

Fundamental Theory of Viscous Incompressible Fluid Flow









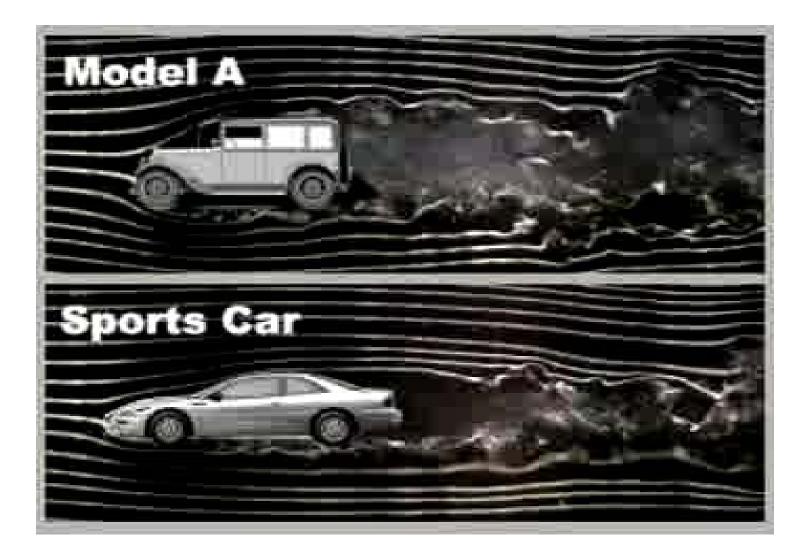




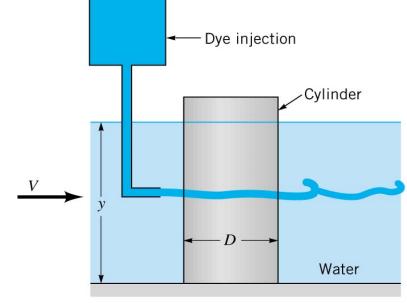


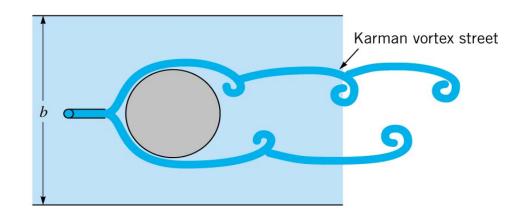




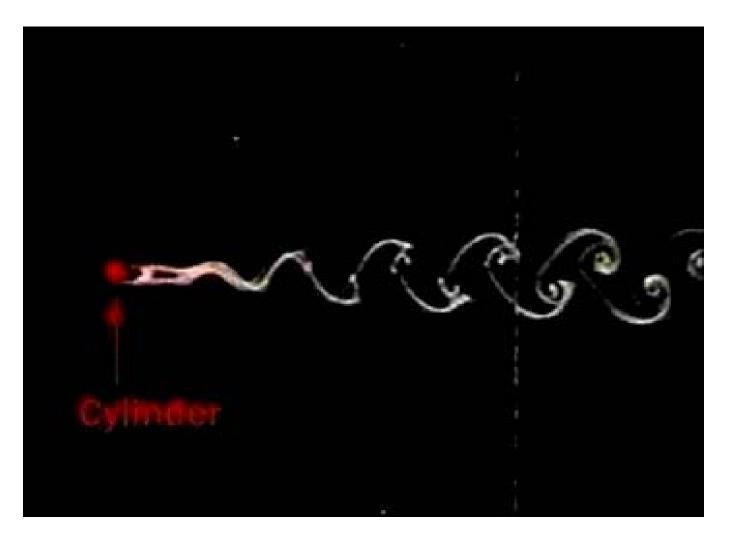


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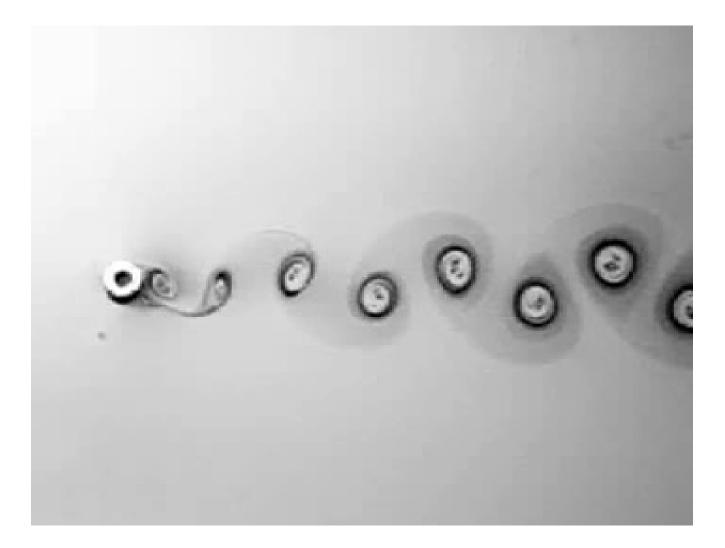




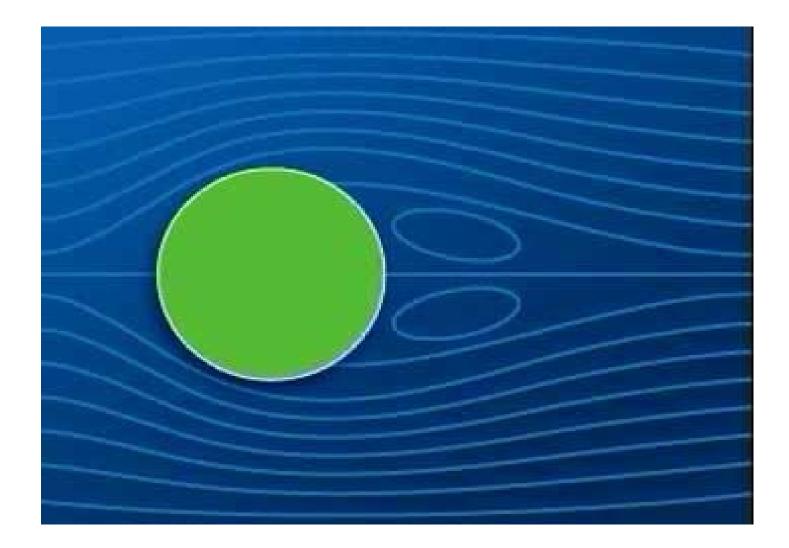




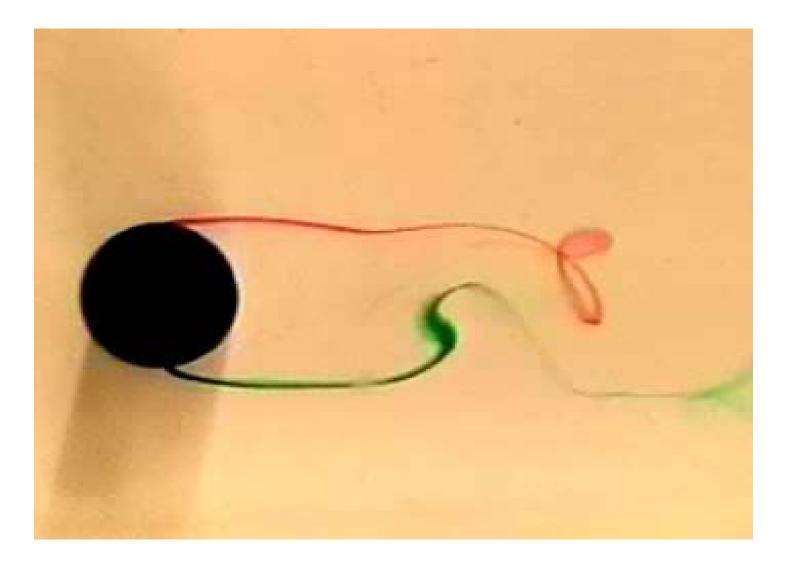




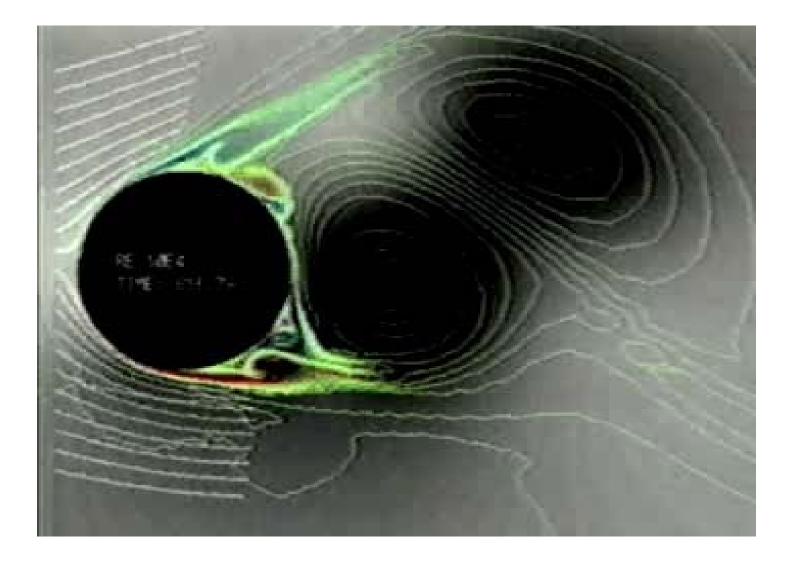




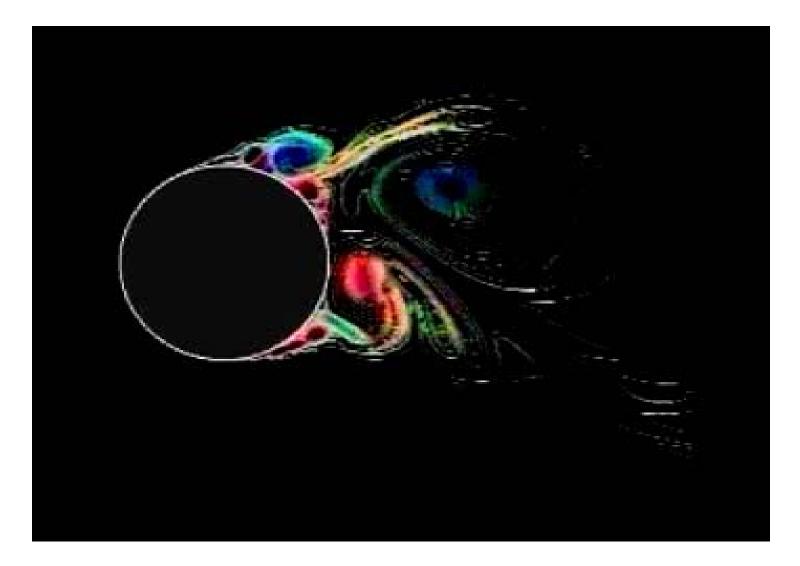




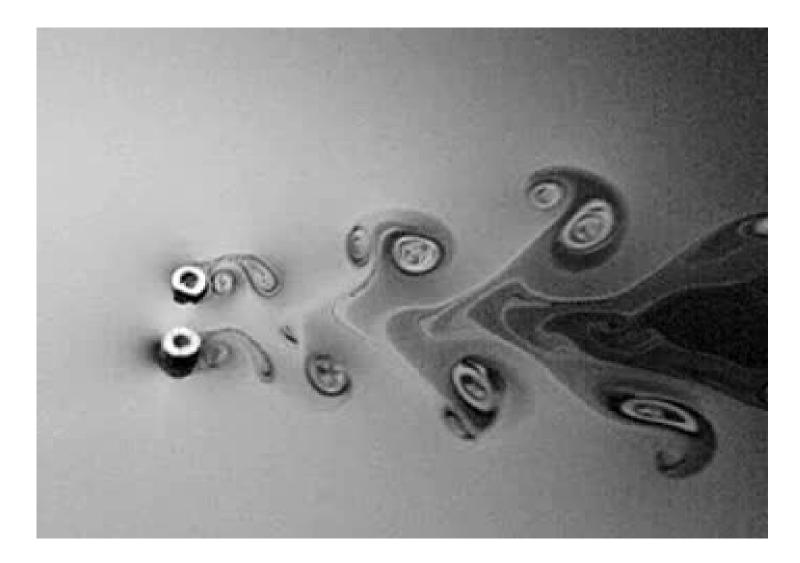




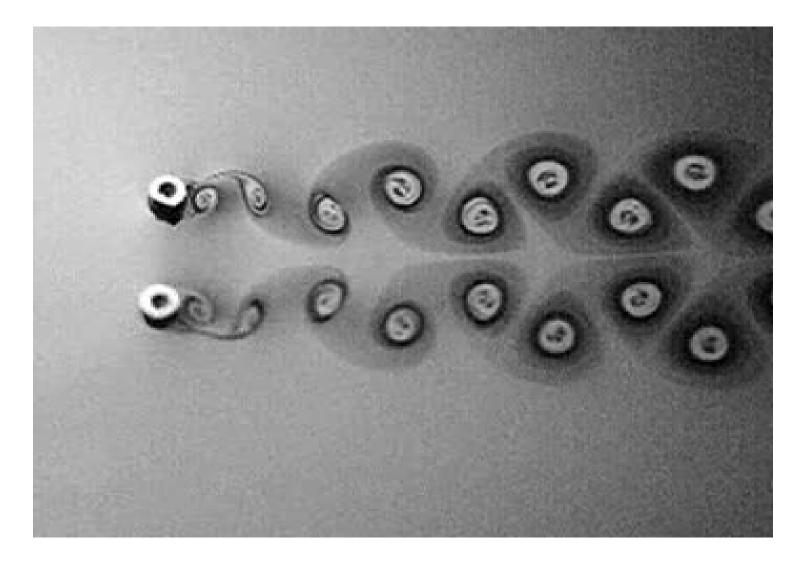




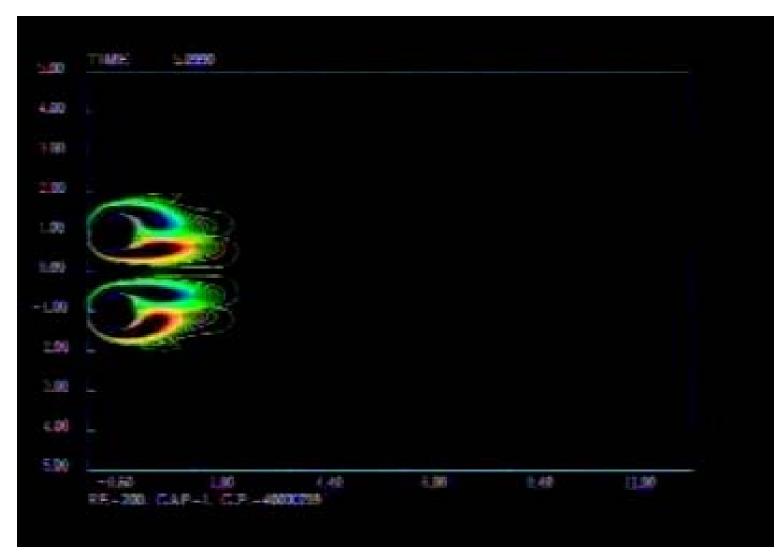




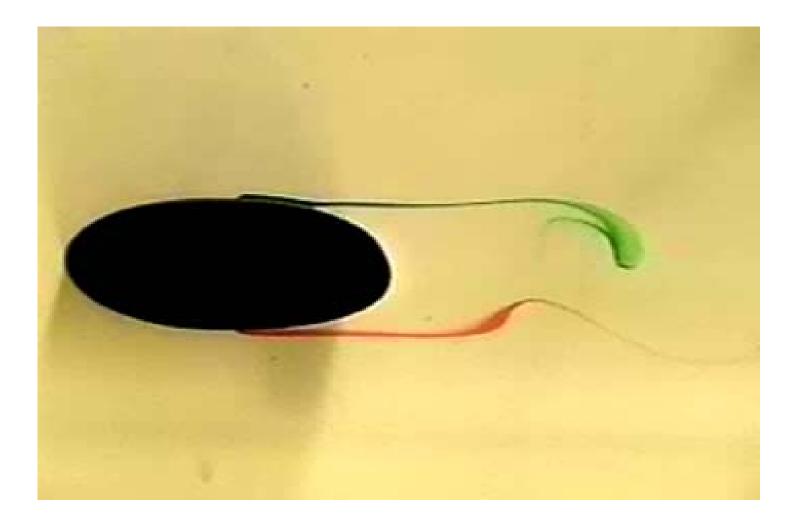




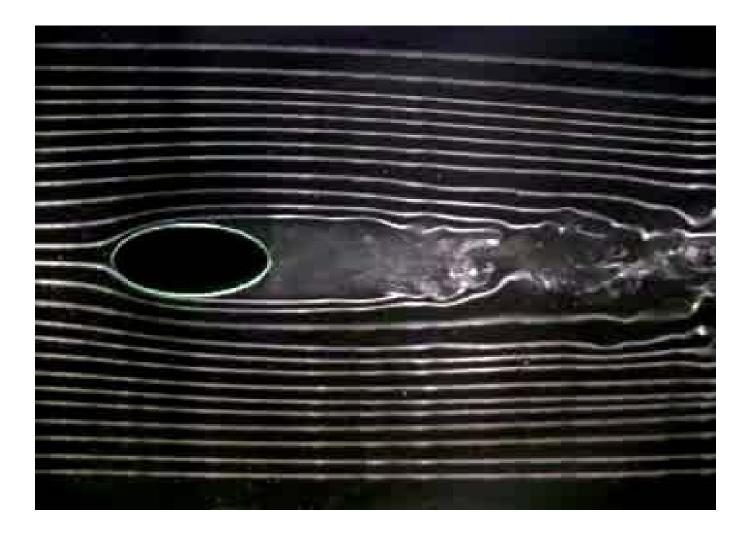




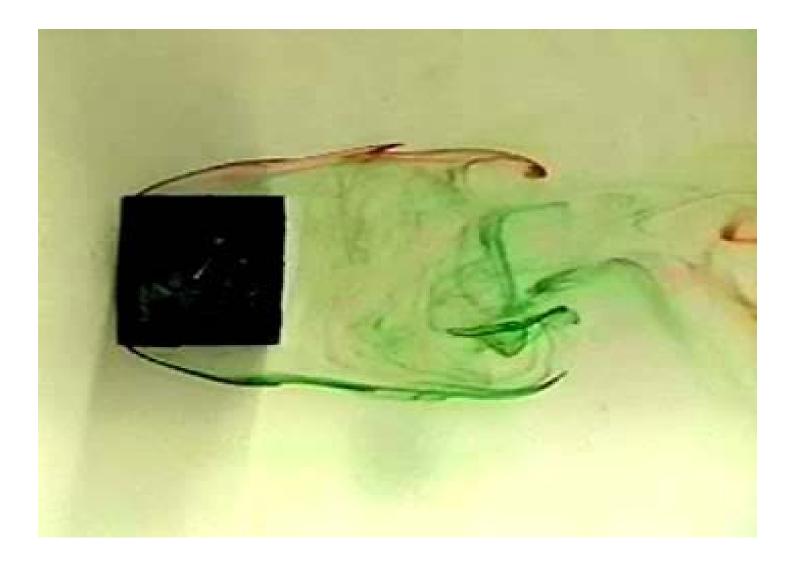




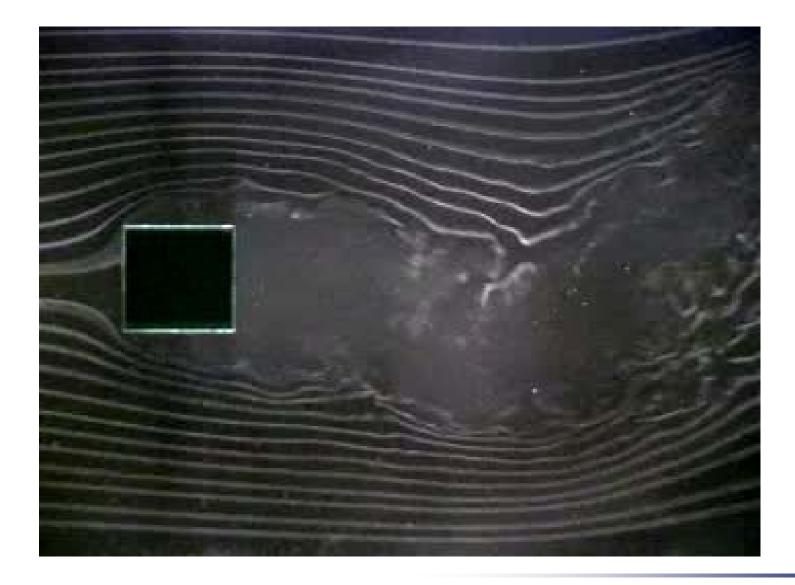




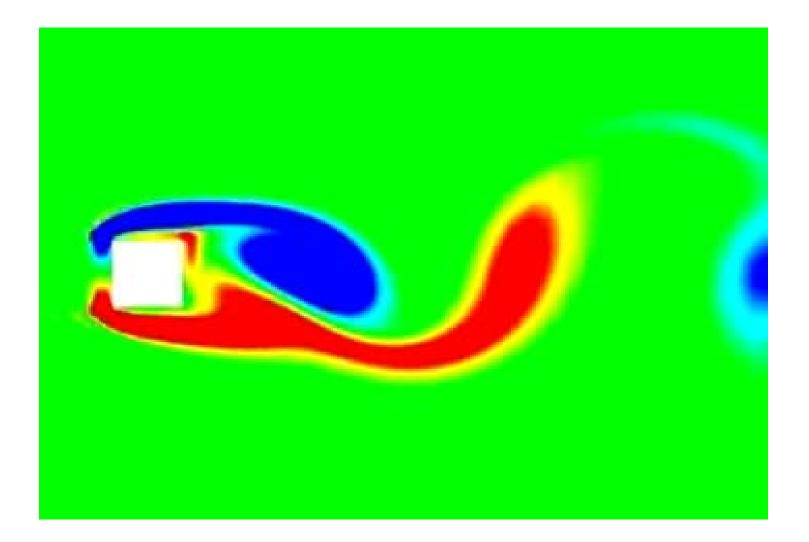




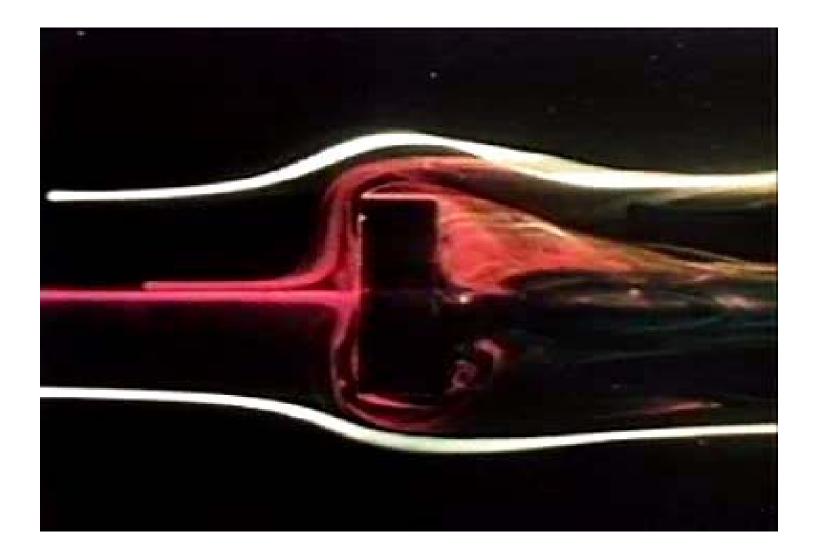




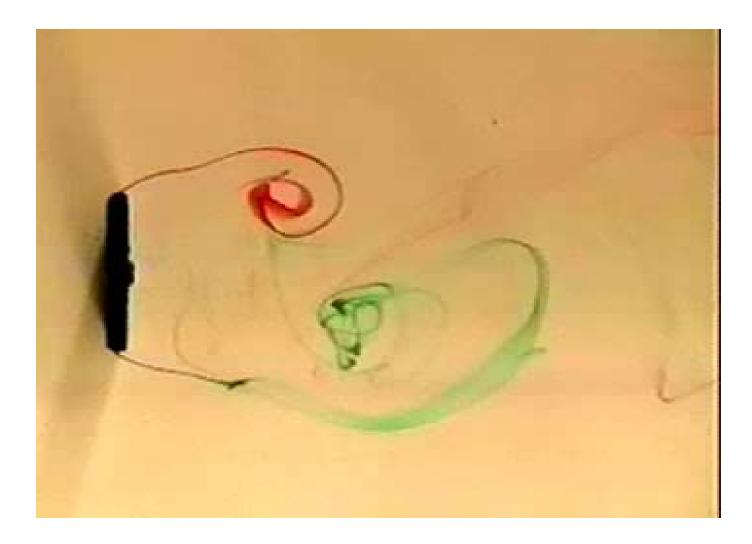




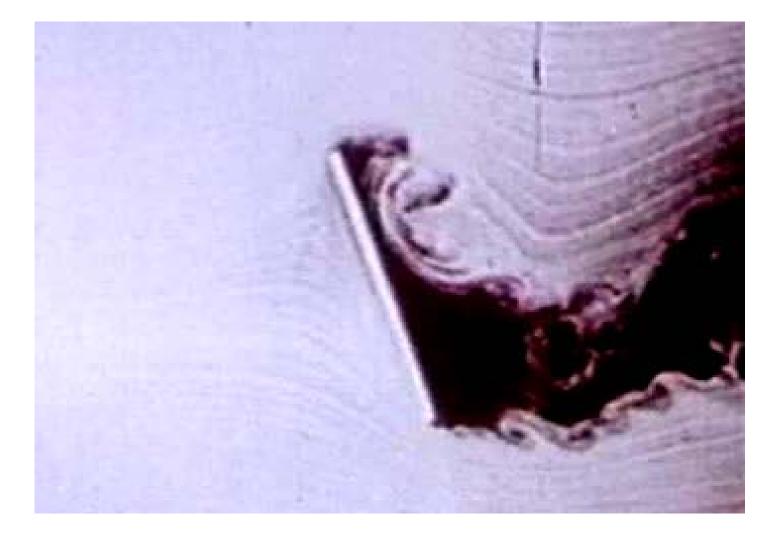




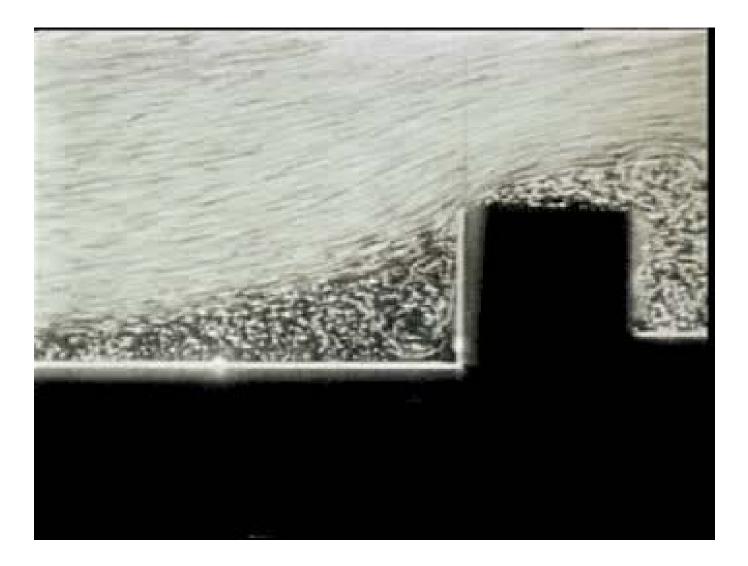




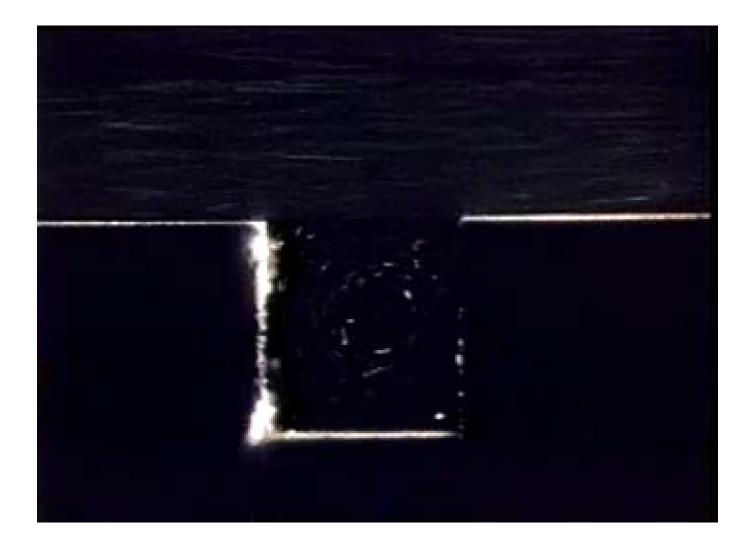




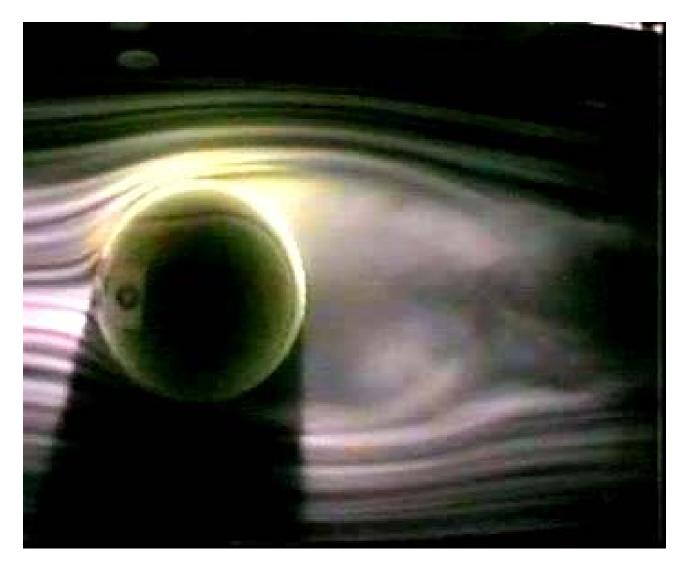




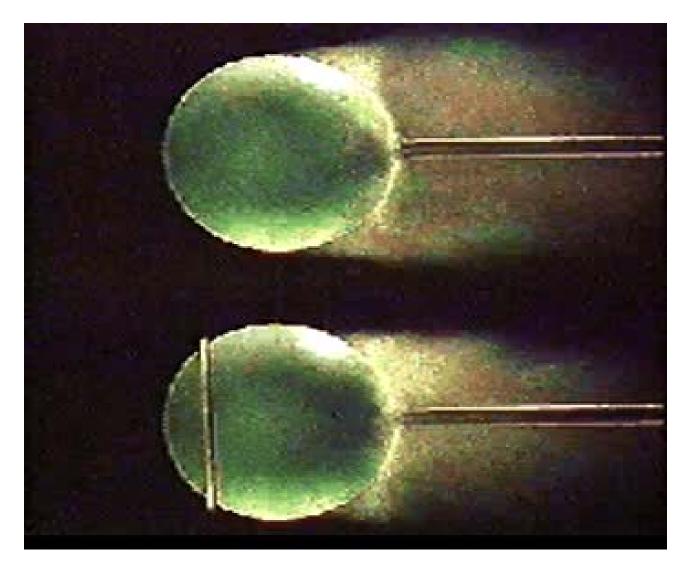
















It can be concluded that viscous flows have characters as follows.

Viscous Flow Phenomena

I) Rotational: vorticity may be non-zero

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- 2) **Dissipation:** mechanical energy may be changed to other energy
- 3) Diffusion: physical quantity may diffuse due to its gradient

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- 4) Unsteady: physical quantity may vary with time
- 5) Unstable: physical quantity may change essentially due to some small disturbance
- 6) Random: physical quantity may change in a random fashion and become indeterminable