



# Introduction to Marine Hydrodynamics (NA235)

Department of Naval Architecture and Ocean Engineering School of Naval Architecture, Ocean & Civil Engineering

Shanghai Jiao Tong University





$$\frac{dx}{u} = \frac{dy}{v} = \frac{dz}{w} = dt$$

# Streamline equation:

$$\frac{dx}{u} = \frac{dy}{v} = \frac{dz}{w}$$



### Motion of fluid elements



Review







#### Shear strain rate – angular deformation

 $\omega_z$  >0, counterclockwise rotation

$$\omega_{x} = \frac{1}{2} \left( \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \right); \quad \omega_{y} = \frac{1}{2} \left( \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right)$$

Rotation with/w.o. deformation?

## 2.3 Deformation and Rotation of Fluid Elements

shear rate tensor 
$$\mathbf{E} = \varepsilon_{ij} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$
  
rate of rotation  $\boldsymbol{\omega} = \omega_x \vec{i} + \omega_y \vec{j} + \omega_z \vec{k} = \frac{1}{2} \boldsymbol{\zeta}$  (vorticity)

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vorticity  $\boldsymbol{\zeta} = \boldsymbol{\Omega} = \nabla \times \mathbf{V}$  (curl of velocity)



## 2.3 Deformation and Rotation of Fluid Elements

Helmholtz decomposition can be written in tensor form:

Let  $\delta \mathbf{r} = dx \mathbf{i} + dy \mathbf{j} + dz \mathbf{k}$ 

 $\mathbf{E} \cdot \delta \mathbf{r} + \boldsymbol{\omega} \times \delta \mathbf{r}$ 

Rate of translation

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Strain rate

Rate of rotation/angular velocity



## **2.4 Rotational and Irrotational Flows**

**Definition:** If the vorticity in a flow field is zero,  $\vec{\Omega}=0$ , the fluid particles are not rotating, the flow in that region is called **irrotational.** Otherwise, the flow in that region is called **rotational.** 

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 $\boldsymbol{\Omega} = \left(\boldsymbol{\Omega}_{x}, \boldsymbol{\Omega}_{y}, \boldsymbol{\Omega}_{z}\right) = 0$ 



## **2.4 Rotational and Irrotational Flows**

## **Application 1**

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The irrotational flow condition in the xy-plane in Cartesian coordinates is:  $\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 0$ 

# Derive the expression of irrotational flow condition in plane polar coordinates.

**Solution**: the relationship between Cartesian coordinates (x, y) and polar coordinates  $(r, \theta)$  is:

$$x = r\cos\theta, \quad y = r\sin\theta$$

where 
$$r = \sqrt{x^2 + y^2}$$
  $\theta = \arctan(\frac{y}{x})$ 

The relationships between the velocities in the two coordinates system are:

$$u = v_r \cos \theta - v_\theta \sin \theta$$
$$v = v_r \sin \theta + v_\theta \cos \theta$$



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The relationships between the differential operators are:

$$\frac{\partial}{\partial r} = \frac{\partial}{\partial x}\frac{\partial x}{\partial r} + \frac{\partial}{\partial y}\frac{\partial y}{\partial r} = \cos\theta\frac{\partial}{\partial x} + \sin\theta\frac{\partial}{\partial y}$$
$$\frac{\partial}{\partial \theta} = \frac{\partial}{\partial x}\frac{\partial x}{\partial \theta} + \frac{\partial}{\partial y}\frac{\partial y}{\partial \theta} = -r\sin\theta\frac{\partial}{\partial x} + r\cos\theta\frac{\partial}{\partial y}$$

Thus:

$$\frac{\partial}{\partial x} = \cos\theta \frac{\partial}{\partial r} - \sin\theta \frac{\partial}{r\partial\theta}$$
$$\frac{\partial}{\partial y} = \sin\theta \frac{\partial}{\partial r} + \cos\theta \frac{\partial}{r\partial\theta}$$

#### The **irrotational flow condition** in Cartesian coordinates is:

$$\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 0$$

**2.4 Rotational and Irrotational Flows** 

Transformed to polar coordinates:

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 $\frac{\partial v}{\partial r} = (\cos\theta \frac{\partial}{\partial r} - \sin\theta \frac{\partial}{r\partial \theta})(v_r \sin\theta + v_\theta \cos\theta)$  $=\sin\theta\cos\theta\frac{\partial v_r}{\partial r} + v_r\cos\theta\cdot 0 - \sin^2\theta\frac{\partial v_r}{r\partial\theta} - \frac{v_r}{r}\sin\theta\cos\theta$  $+\cos^2\theta \frac{\partial v_{\theta}}{\partial r} + v_{\theta}\cos\theta \cdot 0 - \sin\theta\cos\theta \frac{\partial v_{\theta}}{r\partial\theta} + \frac{v_{\theta}}{r}\sin^2\theta$  $\frac{\partial u}{\partial v} = (\sin\theta \frac{\partial}{\partial r} + \cos\theta \frac{\partial}{r\partial \theta})(v_r \cos\theta - v_\theta \sin\theta)$  $=\sin\theta\cos\theta\frac{\partial v_r}{\partial r} + v_r\sin\theta\cdot0 + \cos^2\theta\frac{\partial v_r}{r\partial\theta} - \frac{v_r}{r}\sin\theta\cos\theta$  $-\sin^2\theta \frac{\partial v_\theta}{\partial v_\theta} - v_\theta \sin\theta \cdot 0 - \sin\theta \cos\theta \frac{\partial v_\theta}{\partial v_\theta} - \frac{v_\theta}{\partial v_\theta} \cos^2\theta$ 



 $\Rightarrow$ 

 $\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = (\cos^2 \theta - \sin^2 \theta) \frac{\partial v_{\theta}}{\partial r} + (\sin^2 \theta - \cos^2 \theta) \frac{v_{\theta}}{r} - (\sin^2 \theta + \cos^2 \theta) \frac{\partial v_r}{r \partial \theta} = 0$ 

i.e., 
$$\frac{\partial v_{\theta}}{\partial r} + \frac{v_{\theta}}{r} - \frac{\partial v_{r}}{r\partial \theta} = 0$$

Or, 
$$\frac{\partial (rv_{\theta})}{\partial r} - \frac{\partial v_{r}}{\partial \theta} = 0$$

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## **Application 2**

For a two-dimensional flow field, the velocity is given as:

$$u = -\omega y, v = \omega x$$

Is this flow rotational or irrotational? Will the fluid element deform or

not? Solution:  $\Omega_{z} = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 2\omega, \quad \Omega_{x} = 0, \quad \Omega_{y} = 0$   $\frac{\text{The flow is rotational}}{\varepsilon_{xx}} = \varepsilon_{yy} = \varepsilon_{zz} = 0$   $\varepsilon_{xy} = \varepsilon_{xz} = \varepsilon_{yz} = 0$  There is no deformation of the fluid element. This is used to

describe the flow <u>in</u> the core region of a tornado.

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### **Application 3**

A two-dimensional flow field the velocity is given as: ,

$$u = -\frac{\Gamma}{2\pi} \frac{y}{x^{2} + y^{2}}, \quad v = \frac{\Gamma}{2\pi} \frac{x}{x^{2} + y^{2}}$$

Is this flow rotational or irrotational? Will the fluid element deform or not?

Solution:

$$\omega_{x} = \omega_{y} = 0, \quad \omega_{z} = \frac{1}{2} \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right) = 0 \qquad -$$

$$\varepsilon_{xx} = \frac{\Gamma}{2\pi} \frac{xy}{(x^{2} + y^{2})} = -\varepsilon_{yy} \neq 0, \quad \varepsilon_{zz} \neq 0$$

$$\varepsilon_{xy} = \varepsilon_{yx} = \varepsilon_{zx} = 0$$



The flow is irrotational but deformed, can be used to describe the flow <u>out</u> of the core region of a tornado.

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#### Not all flows with <u>circular</u> streamlines are rotational



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#### **Application 4**



This is a rotational flow with deformation.

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### **Application 5**

Consider a uniform linear motion of a fluid element:

$$u = v_0 = const, \quad v = 0$$

Solution:

$$\varepsilon_{xx} = \varepsilon_{yy} = \varepsilon_{zz} = 0$$
$$\varepsilon_{xy} = \varepsilon_{yz} = \varepsilon_{zx} = 0$$
$$\omega_{x} = \omega_{y} = \omega_{z} = 0$$

This is a irrotational flow without any deformation.



**2.5 Velocity Potential** 

Velocity potential
 <u>for irrotational flow</u>

Stream function
— for incompressible flow

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**2.5 Velocity Potential** 

According to Green's theorem: let P, Q, R be functions of (x, y, z) and have continuous first-order partial derivatives

$$\frac{\partial P}{\partial y}, \frac{\partial P}{\partial z}, \frac{\partial Q}{\partial z}, \frac{\partial Q}{\partial x}, \frac{\partial Q}{\partial x}, \frac{\partial R}{\partial x}, \frac{\partial R}{\partial y}, \cdots$$

in a simple closed curve and its bounded region. If the following equations satisfied in the region:



$$F(x, y, z) = \int P dx + Q dy + R dz$$



**2.5 Velocity Potential** 

The integral is not path-dependent, and has following relationships:

$$\frac{\partial F}{\partial x} = P$$

$$\frac{\partial F}{\partial y} = Q$$

$$\frac{\partial F}{\partial z} = R$$

$$F(x, y, z) = \int \frac{\partial F}{\partial x} dx + \frac{\partial F}{\partial y} dy + \frac{\partial F}{\partial z} dz$$



**2.5 Velocity Potential** 

#### For an *irrotational flow,* $\omega$ =0, thus:



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**2.5 Velocity Potential** 

Thus, there is a function:

$$\phi(x, y, z; t) = \int u dx + v dy + w dz$$

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With the relationships:

$$\frac{\partial \phi}{\partial x} = u$$

$$\frac{\partial \phi}{\partial y} = v$$

$$\frac{\partial \phi}{\partial z} = w$$
i.e,  $\mathbf{V} = \nabla \phi$ 

 $\phi$  is called the velocity potential. A flow has the velocity potential is called **potential flow**.

Irrotational  $\Leftrightarrow \nabla \times V = 0 \Leftrightarrow \phi \Leftrightarrow$  Potential flow





#### Significance of velocity potential:

Velocity potential  $\phi$  is a scalar quantity with one component only:

$$\phi(x, y, .z; t) = \int u dx + v dy + w dz$$

Velocity V is a vector quantity with three components:

$$\frac{\partial \phi}{\partial x} = u$$
$$\frac{\partial \phi}{\partial y} = v$$
$$\frac{\partial \phi}{\partial z} = w$$

Irrotational velocity field  $\Leftrightarrow$  Potential flow



#### **Application**

The velocity field of a rotational flow is given by:

$$u = 2(x-a)y, v = (x+a)^2 - y^2$$

**2.5 Velocity Potential** 

where a is a constant. Consider an irrotational flow, its linear strain rate and shear strain rate are identical to those in rotational flow. The velocity at the origin is 0, i.e., when x = y = 0, u = v = 0. Determine the velocity expression and the velocity potential of this irrotational flow.

**Solution**: for the given rotational flow

linear strain rate: 
$$\varepsilon_{xx} = \frac{\partial u}{\partial x} = 2y, \quad \varepsilon_{yy} = \frac{\partial v}{\partial y} = -2y$$
 (a)  
shear strain rate:  $\varepsilon_{xy} = \frac{1}{2} \left( \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right) = 2x \implies \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} = 4x$  (b)

The linear strain rate and shear strain rate of the irrotational flow are the same, and the irrotational flow has to satisfy:

$$\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 0$$
 (c)

**EXAMPLE VELOCITY Potential**  
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From (c)+(b): 
$$\frac{\partial v}{\partial x} = 2x$$
 i.e.,  $v = x^2 + f(y)$   
Substituting it into (a):  $f'(y) = -2y$ , i.e.,  $f(y) = -y^2 + C_1$   
From (c)-(b):  $\frac{\partial u}{\partial y} = 2x$  i.e.,  $u = 2xy + g(x)$   
Substituting it into (a):  $g'(x) = 0$ , i.e.,  $g(x) = C_2$   
 $C_1$  and  $C_2$  are integral constants. When  $x = y = 0$ ,  $u = v = 0$   
 $\Rightarrow C_1 = C_2 = 0$   
Thus, the velocity is:  $u = 2xy$ ,  $v = x^2 - y^2$   
The velocity potential is:  $\phi = \int u \, dx + v \, dy = \int 2xy \, dx + (x^2 - y^2) \, dy$   
 $= \int_{(0,0)}^{(x,0)} + \int_{(x,0)}^{(x,0)} = 0 + \int_{(x,0)}^{(x,0)} (x^2 - y^2) \, dy = x^2y - \frac{1}{3}y^3$ 

# E 注意文通大学 Review: Key Points in Chapter 2 Shanghai Jiao Tong University

• Two ways of describing a fluid flow:

Lagrangian description, Eulerian description

• Material derivative: 
$$\frac{D()}{Dt} = \frac{\partial()}{\partial t} + (\mathbf{V} \cdot \nabla)()$$

$$\vec{a}_{\text{particle}} = \frac{d\vec{V}_{\text{particle}}}{dt} = \frac{d\vec{V}}{dt} = \frac{d\vec{V}(x_{\text{particle}}, y_{\text{particle}}, z_{\text{particle}}, t)}{dt}$$

$$= \frac{\partial\vec{V}}{\partial t}\frac{dt}{dt} + \frac{\partial\vec{V}}{\partial x_{\text{particle}}}\frac{dx_{\text{particle}}}{dt} + \frac{\partial\vec{V}}{\partial y_{\text{particle}}}\frac{dy_{\text{particle}}}{dt} + \frac{\partial\vec{V}}{\partial z_{\text{particle}}}\frac{dz_{\text{particle}}}{dt}$$

$$\implies \vec{a}_{\text{particle}}(x, y, z, t) = \frac{d\vec{V}}{dt} = \frac{\partial\vec{V}}{\partial t} + u\frac{\partial\vec{V}}{\partial x} + v\frac{\partial\vec{V}}{\partial y} + w\frac{\partial\vec{V}}{\partial z}$$
Eulerian
Eulerian



• Pathline equation:

$$\frac{dx}{u} = \frac{dy}{v} = \frac{dz}{w} = dt$$

• Streamline equation:

dx	$-\frac{dy}{dy}$	dz
U		$\overline{w}$

Pathlines and streamlines are identical in steady flows

#### **Review: Key Points in Chapter 2** :上海交通大學

Motion or deformation of fluid elements 

$$\boldsymbol{\omega} = \frac{1}{2} \boldsymbol{\nabla} \times \mathbf{V} = \omega_x \vec{i} + \omega_y \vec{j} + \omega_z \vec{k} = \left(\frac{\partial w}{\partial y} - \frac{\partial v}{\partial z}\right) \vec{i} + \left(\frac{\partial u}{\partial z} - \frac{\partial w}{\partial x}\right) \vec{j} + \left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}\right) \vec{k}$$

Helmholtz decomposition 

$$\mathbf{V} = \mathbf{V}_0 + \mathbf{E} \cdot \delta \mathbf{r} + \mathbf{\omega} \times \delta \mathbf{r} \qquad \mathbf{E} = \varepsilon_{ij} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

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**Rotational and irrotational flows** 

$$\Omega = \left( \Omega_x, \ \Omega_y, \ \Omega_z \right) = 0$$

**Velocity potential: irrotational flows** 

$$\mathbf{V} = \nabla \phi, \quad \phi(x, y, z; t) = \int u dx + v dy + w dz - \mathbf{w}$$



# Chapter 3 Fluid Dynamics

*Fluid dynamics*: study of the <u>motion</u> of fluids <u>considering</u> the forces and moments that cause the motion.

## **3.1 Representation of Fluid Flow**

Two approaches in fluid dynamics: **system approach** and **control volume approach** 

(1) **System** (material volume) approach: follows the fluid as it moves and deforms; no mass crosses the boundary (Lagrangian description)

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(2) **Control volume** approach: considers the changes in a certain fixed volume; mass can cross the boundary (Eulerian description)



**Reynolds Transport Theorem (RTT)** 

<u>System</u>: a collection of matter of fixed identity (always the <u>same</u> fluid particles), which may move, flow, and interact with its surroundings.

**3.1 Representation of Fluid Flow** 

Material volume: a volume that contains the same fluid as it moves and deforms following the motion of the fluid.

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Material surface: an enclosing surface of a material volume; no fluid particles can cross it.



# <u>Control volume</u>: a volume in space (a geometric entity, independent of mass) through which fluid may flow.

**3.2 Reynolds Transport Theorem** 

**Control volume**: a volume of fluid in a flow field, usually fixed in space, to be <u>occupied</u> by different fluid particles at different times.

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<u>Control surface</u>: imaginary or physical enclosing surface of a control volume, fluid particles can cross it.



# Shanghai Jiao Tong University 3.2 Reynolds Transport Theorem



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## **3.2 Reynolds Transport Theorem**

Control volume at time  $t + \Delta t$ (CV remains fixed in time)

(shaded region)

System (material volume) and control volume at time t

#### **Reynolds Transport Theorem**

Any quantity *G* satisfies the relationship between **material volume** and **control volume**:

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## **3.2 Reynolds Transport Theorem**

For a material volume, after a time period dt:

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$$\left[\iiint_{MV} G(\mathbf{x},t) dV\right]_{t+dt} = \iiint_{MV(t+dt)} G(\mathbf{x},t+dt) dV = \iiint_{MV(t+dt)} \left[G(\mathbf{x},t) + \frac{\partial G}{\partial t} dt + O(dt)^{2}\right] dV$$
The change of material volume in dt is: 
$$\iiint_{MV(t+dt)} \left( \right) = \iiint_{MV} \left( \right) + \iiint_{\Delta V} \left( \right) = \iiint_{CV} \left( \right) + \iiint_{CS} \left( \right) \mathbf{V} \cdot \mathbf{n} dS dt$$
Thus:
$$\left[\iiint_{MV} G(\mathbf{x},t) dV\right]_{t+dt} = \iiint_{MV(t+dt)} \left[G(\mathbf{x},t) + \frac{\partial G}{\partial t} dt\right] dV$$

$$= \iiint_{CV} \left[G(\mathbf{x},t) + \frac{\partial G}{\partial t} dt\right] dV + \iint_{CS} \left[G(\mathbf{x},t) + \frac{\partial G}{\partial t} dt\right] \mathbf{V} \cdot \mathbf{n} dS dt$$

$$= \iiint_{CV} G(\mathbf{x},t) dV + \left[\iiint_{CV} \frac{\partial G}{\partial t} dV + \iint_{CS} G(\mathbf{x},t) \mathbf{V} \cdot \mathbf{n} dS \right] dt$$
RTT can be derived from the equation above:
$$\frac{d}{dt} \iiint_{MV} GdV = \left\{\left[\iiint_{MV} GdV\right]_{t+dt} - \iiint_{MV} GdV\right] / dt = \iiint_{CV} \frac{\partial G}{\partial t} dV + \iint_{CS} GV \cdot \mathbf{n} dS \text{ Control surface (CS)} \right\}$$





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# **3.3 Continuity Equation**

**Continuity Equation:** conservation of mass (mass of a system is neither be created nor destroyed)

In RTT equation, if the quantity is mass, i.e.,  $G = \rho$ , then:

L.H.S. 
$$\frac{d}{dt} \iiint_{MV} \rho d\Psi = \frac{d}{dt} (\text{mass in } MV) = 0$$

(From the definition of MV: it always contains the same fluids)

R.H.S. 
$$\frac{\partial}{\partial t} \iiint_{CV} \rho d \Psi + \iint_{CS} \rho V \cdot n d A$$
$$= \underbrace{\iiint_{CV} \frac{\partial \rho}{\partial t} d\Psi}_{CV \text{ is stationary}} + \underbrace{\iiint_{CV} \nabla \cdot (\rho V) d\Psi}_{\text{by Gauss theorem}}$$